

## The Crystal Structure of $\text{KCeS}_2$

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(Received 28 May 1975; accepted 3 December 1975)

The crystal structure of  $\text{KCeS}_2$  has been determined by single-crystal X-ray diffraction techniques. The compound crystallizes with the rhombohedral  $\alpha\text{-NaFeO}_2$  structure,  $a=4.228$ ,  $c=21.80$  Å (hexagonal setting). Full-matrix least-squares refinement, assuming  $R\bar{3}m$  symmetry, of 307 observed symmetry-independent three-dimensional diffractometer data (Mo  $K\alpha$ ) gave a conventional  $R$  index of 0.023 ( $R_{wF}=0.024$ ). Both cations are coordinated octahedrally; the possibility of a trigonal prismatic coordination of one of the cations was rejected on the basis of the refinement results. Refinement under  $R3m$  symmetry yielded no significant improvement of  $R_{wF}$ . Under the assumption that  $R\bar{3}m$  is the right space group, the structure can be regarded as a layer-ordered NaCl structure, in which the close-packed sulphur layers are displaced towards the cerium layers.

### Introduction

Since 1964, compounds of the type  $\text{MLnS}_2$  ( $M=\text{Li, Na, K}$ ;  $\text{Ln}=\text{rare earth ion}$ ) have been known. Ballestracci & Bertaut (1964) and Ballestracci (1965) determined their structures by X-ray powder measurements. They reported these compounds to crystallize with the  $\alpha\text{-NaFeO}_2$  structure, space group  $R\bar{3}m$ , or the NaCl structure.

For several reasons, however, it seems important to perform a single-crystal X-ray study of these sulphides. The ascertainment of the space group from powder data is ambiguous. It is important to establish definitely whether the coordination of the cations is octahedral or trigonal prismatic (Jellinek, 1960). For the interpretation of the physical properties of these compounds, e.g. crystal field splitting, spectra and magnetic properties at low temperatures, the exact value of the positional parameter of the sulphur ions is of interest. Finally, the crystal chemistry of the  $\text{MLnS}_2$  compounds has not been studied in detail. Bronger, Elter, Maus & Schmitt (1973) found the rubidium and some of the caesium compounds to have the  $\alpha\text{-NaFeO}_2$  structure too; so the same structure would occur for a wide range of the radius ratio of the cations ( $0.62 < r_{\text{Ln}^{3+}}^*/r_{\text{S}^{2-}}^* < 1.23$ ).

### Experimental

Crystals of  $\text{KCeS}_2$  were prepared by heating a mixture of  $\text{CeO}_2$  (99.9%, Alfa Inorganics) and an excess of  $\text{K}_2\text{S}$  (p.a. Merck) in a stream of  $\text{H}_2\text{S}$ . After cooling from  $1000^\circ\text{C}$  ( $30^\circ/\text{h}$ ) the product was washed with water and acetone, and dried. Unit-cell dimensions, derived from powder diffractograms, are  $a=4.228$  (4),  $c=21.80$  (1) Å (hexagonal axes). The calculated density for  $Z=3$  and a molecular weight of 243.4,  $d=3.59$   $\text{g cm}^{-3}$ , agrees with the measured density,  $d=3.61$   $\text{g cm}^{-3}$ . A hexagonal-shaped crystal of approximate size  $0.1 \times 0.1 \times 0.05$  mm was mounted on an En-

raf-Nonius three-circle single-crystal diffractometer, with the hexagonal  $a$  axis along the  $\phi$  axis.

Intensities were recorded by the  $\theta$ - $2\theta$  scan method for all reflexions with  $\theta$  between  $5$  and  $40^\circ$ , using graphite monochromatized Mo  $K\alpha$  radiation. Background intensities were determined at  $\theta \pm \frac{1}{2}\Delta$ , where  $\Delta=0.8+1.2 \tan \theta$ . The mean counting time was 36 s for each background and 72 s for the scan. The total number of reflexions measured by this method was 1740. Reflexions with intensities less than twice the standard deviation, calculated from counting statistics, were considered as not significant. All data were corrected for Lorentz and polarization effects and for absorption (de Graaff, 1973), using transmission factors between 0.13 and 0.41 ( $\mu=118.0$   $\text{cm}^{-1}$ ).

All intensities were reduced to  $F$  values and an averaging procedure was carried out, giving a total number of 307 observed symmetry-independent reflexions. Starting values for the scaling factor and the initial overall isotropic thermal parameter  $B$  were obtained from a calculated Wilson plot.

### Refinement of the structure

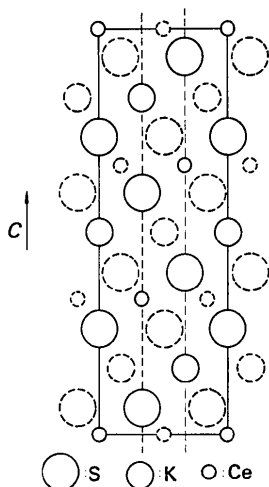
The function minimized during the least-squares refinement process was  $\sum w_F(|F_o| - |F_c|)^2$ , with the weighting factor  $w_F=(\sigma_F)^{-2}$ . Discrepancy indices referred to are:  $R_F = \sum ||F_o| - |F_c|| / \sum |F_o|$  and  $R_{wF} = [\sum w_F(|F_o| - |F_c|)^2 / \sum w_F |F_o|^2]^{1/2}$ .

From zero- and upper-level Weissenberg photographs the crystal symmetry and the approximate cell parameters were determined. Systematically absent reflexions were of the type  $-h+k+l \neq 3n$ , indicating a rhombohedral symmetry, and the intensities of the reflexions  $hkil$  and  $khil$  were equal. The Laue group appeared to be  $\bar{3}m$ , leaving three possible space groups:  $R\bar{3}m$ ,  $R3m$  and  $R32$ .

Two models with both cations in an octahedral coordination and two models with either Ce or K in a

Table 1. Positions for the atoms in the different space groups

	Positions	Space group	Cation coordination
Ce	3(a) 0,0,0	$R\bar{3}m$ and $R32$	both octahedral
K	3(b) 0,0, $\frac{1}{2}$		
S	6(c) 0,0,z } $z \approx \frac{1}{4}$ 0,0, $\bar{z}$ }		
Ce	3(a) 0,0,z; z=0	$R3m$	both octahedral
K	3(a) 0,0,z; z $\approx \frac{1}{2}$		
S(1)	3(a) 0,0,z; z $\approx \frac{1}{4}$		
S(2)	3(a) 0,0,z; z $\approx \frac{3}{4}$		
Ce	3(a) 0,0,z; z=0	$R3m$	Ce octahedral
K	3(a) 0,0,z; z $\approx \frac{5}{8}$		K prismatic
S(1)	3(a) 0,0,z; z $\approx \frac{5}{12}$		
S(2)	3(a) 0,0,z; z $\approx \frac{7}{12}$		
Ce	3(a) 0,0,z; z $\approx \frac{5}{8}$	$R3m$	Ce prismatic
K	3(a) 0,0,z; z=0		K octahedral
S(1)	3(a) 0,0,z; z $\approx \frac{5}{12}$		
S(2)	3(a) 0,0,z; z $\approx \frac{7}{12}$		

Fig. 1. Projection on  $(11\bar{2}0)$  of  $KCeS_2$ .

trigonal prismatic coordination, were refined by a full-matrix least-squares method. The positions for the atoms in the different space groups, according to these models, are given in Table 1. Refinement of both prismatic models (with either Ce or K on  $z \approx \frac{5}{8}$ ) gave poor results: the best  $R_F$  value obtained was 0.40.

For the octahedral model, however, refinement under  $R\bar{3}m$  or  $R3m$  symmetry of only the scaling factor and the overall thermal parameter yielded  $R_F=0.17$ . Three cycles of full-matrix refinement under  $R\bar{3}m$  symmetry of the positional parameter of sulphur and the isotropic thermal parameters of all atoms reduced  $R_F$  to 0.04. At this stage the thermal parameters were allowed to refine anisotropically and an extinction correction was carried out. This led to final discrepancy indices, considering  $R\bar{3}m$  symmetry,  $R_F=0.023$  and  $R_{wF}=0.024$ . Having reached this level, it can be stated already that  $R\bar{3}m$  is the most likely space group. The low indices  $R_F$  and  $R_{wF}$ , the low value of the parameter  $U_{33}$  of sulphur and the absence of chemical arguments for lowering the symmetry strongly indicate this. Still it was thought useful to attempt a refinement under  $R3m$  symmetry, decoupling the sulphur positions and allowing the  $z$  parameter of K to refine also, with the expectation that adding new parameters to be refined would reduce the discrepancy indices to a lower level. We carried out three refinements under  $R3m$  symmetry. First the K position was kept exactly at  $(0,0,\frac{1}{2})$  and the displacements of S(1) and S(2) were restricted to be equal and in the same direction for both sulphur ions ( $\Delta z_{S(1)} = \Delta z_{S(2)}$ ). Then, in the second refinement the  $z$  parameter of K was allowed to refine also. Finally, the  $z$  parameters of K(1), S(1) and S(2) were all refined independently. In all cases the anisotropic thermal parameters of all atoms were refined as well, the starting values being the final values from the refinement under  $R\bar{3}m$ . The thermal parameters for the sulphur atoms 1 and 2 were restricted to keep equal values for both atoms. The results are summarized in Table 2.

Table 2. Positional and thermal parameters ( $\times 10^4$ ) for  $KCeS_2$ 

The general anisotropic temperature factor has the form  $\exp[-2\pi^2(\sum \sum U_{ij}h_i h_j a_i^* a_j^*)]$ . Owing to the site symmetry of the ions there are several restrictions on the thermal parameters  $U_{ij}$ :  $U_{11} = U_{22} = 2U_{12}$ ;  $U_{23} = U_{31} = 0$ . Estimated standard deviations in the least significant digits are in parentheses.

Space group		x	y	z	$U_{11} = U_{22} = 2U_{12}$	$U_{33}$	$R_F(\%)$	$R_{wF}(\%)$
$R\bar{3}m$	Ce	0	0	0	69 (1)	94 (1)	2.29	2.36
	K	0	0	5000	181 (3)	134 (4)		
	S	0	0	2634.6 (4)	105 (2)	103 (3)		
$R3m(1)$	Ce	0	0	0	69 (1)	95 (1)	2.30	2.29
	K	0	0	5000	180 (3)	130 (4)		
	S(1)	0	0	2652 (3)	99 (2)	78 (7)		
	S(2)	0	0	7382 (3)	99 (2)	78 (7)		
$R3m(2)$	Ce	0	0	0	69 (1)	95 (1)	2.30	2.28
	K	0	0	5010 (10)	181 (3)	121 (12)		
	S(1)	0	0	2647 (6)	99 (2)	84 (8)		
	S(2)	0	0	7377 (6)	99 (2)	84 (8)		
$R3m(3)$	Ce	0	0	0	69 (1)	95 (1)	2.30	2.28
	K	0	0	5018 (15)	181 (3)	110 (25)		
	S(1)	0	0	2643 (7)	99 (2)	88 (8)		
	S(2)	0	0	7374 (7)	99 (2)	88 (8)		

As can be seen, refinement under *R3m* gives a slight improvement of the index *R*<sub>wF</sub>. However, this improvement is not very significant because of the almost complete correlation of the sulphur parameters, *U*<sub>33</sub> and *z*, which are refined, as is shown by the correlation matrices. So we conclude that *R3m* is the most likely space group for KCeS<sub>2</sub>. The observed and calculated structure factors in that space group are listed in Table 3.

Table 3. Structure factors

h	k	l	F <sub>obs</sub>	F <sub>calc</sub>	h	k	l	F <sub>obs</sub>	F <sub>calc</sub>	h	k	l	F <sub>obs</sub>	F <sub>calc</sub>	h	k	l	F <sub>obs</sub>	F <sub>calc</sub>	h	k	l	F <sub>obs</sub>	F <sub>calc</sub>
0	0	0	4520	4517	1	0	12	3251	3081	2	0	-4	5677	5977	3	-1	16	3133	3056	4	-2	3	2266	2258
0	0	6	4253	3865	1	0	18	3186	3053	2	0	-1	2912	2906	3	-1	19	2912	2906	4	-2	6	2076	2156
0	0	12	1601	1424	1	0	25	727	734	2	0	0	2845	2944	3	-1	22	2953	2522	4	-2	9	1155	1147
0	0	18	3538	3380	1	0	31	1546	1551	2	0	1	1915	1923	3	-1	25	565	576	4	-2	12	2090	2186
0	0	24	3857	3657	1	0	31	1600	1460	2	0	8	5117	5247	3	-1	28	1301	1331	4	-2	15	2517	2577
0	0	30	3538	3380	1	0	34	2037	2051	2	0	11	3395	3339	3	-1	31	1261	1281	4	-2	18	2565	2532
0	0	36	4921	4746	1	0	37	825	851	2	0	14	3056	2992	3	-1	34	1765	1785	4	-2	21	452	451
0	0	42	2262	2185	1	0	37	876	876	2	0	17	4012	3952	3	-1	37	706	743	4	-2	24	1584	1558
0	0	48	2038	2027	1	0	37	3077	3171	2	0	20	2716	2597	3	-1	40	1800	1821	4	-2	27	1438	1514
0	0	54	850	846	1	0	37	3215	3166	2	0	23	2627	2551	3	-1	43	1845	1881	4	-2	30	1480	1514
0	0	60	961	946	1	0	37	1477	1419	2	0	26	2627	2551	3	-1	46	1845	1881	4	-2	33	1480	1514
0	0	66	438	428	1	0	37	1546	1551	2	0	29	1890	1890	3	-1	49	1845	1881	4	-2	36	1480	1514
0	0	72	1672	1686	1	0	37	1546	1551	2	0	32	1890	1890	3	-1	52	1845	1881	4	-2	39	1480	1514
0	0	78	1018	1035	1	0	37	1546	1551	2	0	35	1890	1890	3	-1	55	1845	1881	4	-2	42	1480	1514
0	0	84	1150	1186	1	0	37	1546	1551	2	0	38	1890	1890	3	-1	58	1845	1881	4	-2	45	1480	1514
0	0	90	778	784	1	0	37	1546	1551	2	0	41	1890	1890	3	-1	61	1845	1881	4	-2	48	1480	1514
0	0	96	2655	2601	1	0	37	1546	1551	2	0	44	1890	1890	3	-1	64	1845	1881	4	-2	51	1480	1514
0	0	102	3533	3530	1	0	37	1546	1551	2	0	47	1890	1890	3	-1	67	1845	1881	4	-2	54	1480	1514
0	0	108	3005	2889	1	0	37	1546	1551	2	0	50	1890	1890	3	-1	70	1845	1881	4	-2	57	1480	1514
0	0	114	3568	3369	1	0	37	1546	1551	2	0	53	1890	1890	3	-1	73	1845	1881	4	-2	60	1480	1514
0	0	120	651	631	1	0	37	1546	1551	2	0	56	1890	1890	3	-1	76	1845	1881	4	-2	63	1480	1514
0	0	126	3768	3768	1	0	37	1546	1551	2	0	59	1890	1890	3	-1	79	1845	1881	4	-2	66	1480	1514
0	0	132	5718	6459	1	0	37	1546	1551	2	0	62	1890	1890	3	-1	82	1845	1881	4	-2	69	1480	1514
0	0	138	4145	3888	1	0	37	1546	1551	2	0	65	1890	1890	3	-1	85	1845	1881	4	-2	72	1480	1514
0	0	144	3719	3523	1	0	37	1546	1551	2	0	68	1890	1890	3	-1	88	1845	1881	4	-2	75	1480	1514
0	0	150	1126	1056	1	0	37	1546	1551	2	0	71	1890	1890	3	-1	91	1845	1881	4	-2	78	1480	1514
0	0	156	4761	3977	1	0	37	1546	1551	2	0	74	1890	1890	3	-1	94	1845	1881	4	-2	81	1480	1514

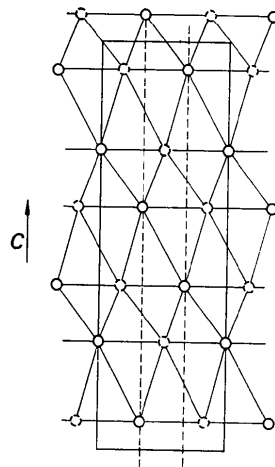


Fig. 2. Projection on (11 $\bar{2}$ 0), showing the shift of the sulphur layers.

The size of the distortion from the ideal structure can also be described by a position parameter  $\epsilon$ , indicating the shift of the sulphur ions (Brunel, de Bergevin & Gondrand, 1972). From the determined *z* parameter for the sulphur ion in KCeS<sub>2</sub> we obtain  $\epsilon_{\text{obs}} = -0.0135$  (9). Considering the ions as hard spheres one can calculate  $\epsilon$  according to the following formula:

$$\epsilon = \frac{1}{4} - \frac{(r_{\text{Ce}}^{3+} + r_{\text{S}}^{2-})}{2(r_{\text{Ce}}^{3+} + r_{\text{K}}^+ + 2r_{\text{S}}^{2-})}$$

(in units of the hexagonal *c* axis). Using the ionic radii given by Shannon & Prewitt (1969) we obtain  $\epsilon_{\text{calc}} = -0.0148$ , which is in fairly good agreement with the observed value.

A remaining question is what kind of ordering will occur in the case where the alkali ion is much smaller than the rare-earth ion. An attempt to prepare single crystals of some Li compounds is in progress in order to solve this problem.

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Discussion

The structure can be described as an ordering in the cation positions in the NaCl structure, or as the MgCl<sub>2</sub> structure in which the Mg positions are occupied by the Ce ions and the vacant octahedral positions by the K ions. The projection of the structure on (11 $\bar{2}$ 0) is given in Fig. 1.

The regular S<sub>6</sub> octahedra of the idealized structure are distorted: the CeS<sub>6</sub> octahedron is compressed (reduced *c/a* ratio is 0.7205) and the KS<sub>6</sub> octahedron is elongated (reduced *c/a* 0.9982) along the threefold axis (*cf.* Fig. 2). As could be expected, the octahedra with the highest charged cations are the least distorted ones.